

Development of a 3-D, multi-nuclear continuous wave NMR imaging system

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Abstract

The development of a 3-D, multi-nuclear continuous wave NMR imaging (CW-NMRI) system is described and its imaging capability is demonstrated on a range of materials exhibiting extremely short T_2 relaxation values. A variety of radiofrequency resonators were constructed and incorporated into a new gradient and field offset coil assembly, while the overall system design was modified to minimise microphonic noise which was present in an earlier prototype system. The chemically combined ^{27}Al in a high temperature refractory cement was imaged, and the CW-NMRI system was found to be sensitive to small differences in ^{27}Al content in these samples. The penetration of ^{23}Na in salt water into samples of ordinary Portland cement (OPC) was investigated, with enhanced uptake observed for samples with larger pore size distributions. The solid ^{13}C component in a carbonated cement sample was also imaged, as were the ^7Li nuclei in a sample of powdered Li_2CO_3 . A spatial resolution of 1 mm was measured in an image of a rigid polymeric material exhibiting a principal T_2^* value of 16.3 μs . Finally, a high-resolution 3-D image of this rigid polymer is presented. © 2005 Elsevier Inc. All rights reserved.

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1. Introduction

The interest in imaging solid materials and liquids in confined geometries has led to the ongoing development of a range of solid state NMR imaging techniques over the past 20 years. Such systems present extreme difficulties for conventional NMR imaging due to the short T_2 relaxation values and broad NMR linewidths which these materials exhibit. These problems can be overcome, to some degree, by decreasing the effective linewidth of the sample using techniques which reduce the line-broadening mechanisms within the sample, so that

lower strength field gradients may be used (line-narrowing techniques). This can be achieved by using a range of special multi-pulse radiofrequency (RF) irradiation methods [1–6] or by physically spinning the sample at the magic angle [7–10]. Another approach to solid state NMR imaging which has been exploited in several techniques involves the application of sufficiently large magnetic field gradients to produce spatial localisation on the resonant spins, thereby retaining the broad linewidths which contain useful information from the T_2 relaxation phenomena which broaden the lines. Examples include the stray field imaging (STRAFI) technique, which typically uses the extremely strong static field gradients present in the fringe field of solenoidal superconducting magnets [11,12], and the oscillating gradient technique, which produces large, sinusoidally varying gradients by incorporating the gradient coils into a

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resonant circuit [13]. A further approach to solid state NMR imaging, which uses moderately strong gradients, is single-point imaging (SPI, or its variant SPRITE, single-point ramped imaging with T_1 enhancement). This technique uses pure phase encoding and effectively eliminates all line-broadening mechanisms by acquiring a single data sample at a fixed time after the application of an RF pulse, which is repeated as the gradient is incremented [14–16].

Despite the wealth of techniques that have been developed to image materials in the solid state, no single technique has emerged as an all round gold standard, with each technique demonstrating certain advantages and disadvantages compared to the others. For example, line-narrowing techniques often require the application of high RF power and place severe restrictions on sample size. Together with the complexity of the experimental set-up and the demanding equipment requirements, their application to the imaging of solid materials has been limited [17]. The unidirectional nature of the strong gradient in the STRAFI experiment makes it difficult and time consuming to acquire images in more than one dimension since both sample translation and rotation are required, and consequently STRAFI has found application predominately in the imaging of materials where the process of interest has been reduced to 1-D, for example, profiling perpendicular to thin films, where high-resolution imaging of the in-plane processes is not so critical [18,19]. Nevertheless, some 3-D studies have been reported (for example, see [12,20]). The SPI/SPRITE technique, while perhaps the most successful technique to date from the point of view of its range of applications reported in the literature, is limited by equipment deadtime which imposes a lower limit on the shortest T_2 value which can be studied [21,22].

A newly developed approach to imaging materials in the solid state is continuous wave NMR imaging (CW-NMRI), a frequency encoding technique which uses continuous RF irradiation and detection in the presence of continuously applied, moderately strong gradients [23–25]. Extremely short ($<10 \mu\text{s}$) T_2 materials are amenable to study using this technique due to the lack of experimental deadtime, although imaging using a prototype system have thus far been limited to one and two dimensions [25]. Its extension to 3-D imaging is described in this paper.

Given the difficulties that exist with imaging ^1H in the solid state, it is perhaps not surprising that there are considerably fewer reported studies in the literature which describe the imaging of other nuclei in the solid state. Nevertheless, the ability to image such nuclei raises a host of interesting possible applications, for example, the imaging of ^{31}P in bone and ^{27}Al in clay minerals and cementitious materials. STRAFI has been used to produce echo trains and 1-D line profiles from a wide variety of nuclei in various test compounds and

phantoms [26–29]. However, no 2- or 3-D non-proton STRAFI images have been reported in the literature. SPRITE has been used to image liquid-phase ^{23}Na and ^{35}Cl in salt water penetrating into a cement, but again only 1-D profiles were reported [30]. However, Halse et al. [31] recently succeeded in obtaining a 2-D ^{19}F image of a PTFE object. The second major development to the CW-NMRI technique described in this paper is the addition of multi-nuclear imaging capabilities, and examples of the imaging of solid phases of ^7Li , ^{13}C , ^{23}Na , and ^{27}Al in a variety of material systems are presented. Coupled with its 3-D imaging capability, these developments pave the way for a broad range of new applications for the CW-NMRI technique.

2. Experimental

2.1. CW-NMRI basics

In continuous wave NMRI, a sample is exposed to continuous RF irradiation at one specific frequency and its absorption signal is measured by sweeping the applied magnetic field through resonance. Operation of our spectrometer is described with reference to the block diagram in Fig. 1. The RF, supplied by a synthesiser (Hewlett–Packard, USA, Model HP8647A), is applied to one terminal of a hybrid junction (Lorch Electronics, USA, Model JH280E). The two neighbouring ports of the hybrid junction are connected to a 50Ω load and the resonator, which is matched to 50Ω , while the signal appearing on the opposite port is sent to a home-built diode detector. When off-resonance, power from the synthesiser is divided equally between the 50Ω load and the resonator. Sweeping of the magnetic field via a solenoidal field offset coil then brings the spins in the sample into resonance, which results in a change in impedance of the resonator. The resulting impedance mismatch across the hybrid junction causes power to be transmitted through the junction into the detector. To improve signal sensitivity, a small-amplitude audiofrequency modulation ($\sim 1 \text{ kHz}$) is superimposed on the swept magnetic field, and lock-in detection is used to extract the signal with the same frequency as the modulation and at a fixed relative phase. This phase-sensitive detection allows for the use of an extremely narrow bandwidth detector, resulting in a large increase in the signal-to-noise ratio (SNR) of the measured signal. The output from the lock-in amplifier (Stanford Research Systems, USA, Model 830 DSP) is proportional to the change in signal reflected from the resonator as the field is swept, and thus the measured first derivative signal must be integrated to recover the sample's absorption curve.

Imaging with the CW-NMRI technique is achieved by frequency encoding in one direction only, using a moderately strong, static magnetic field gradient. The

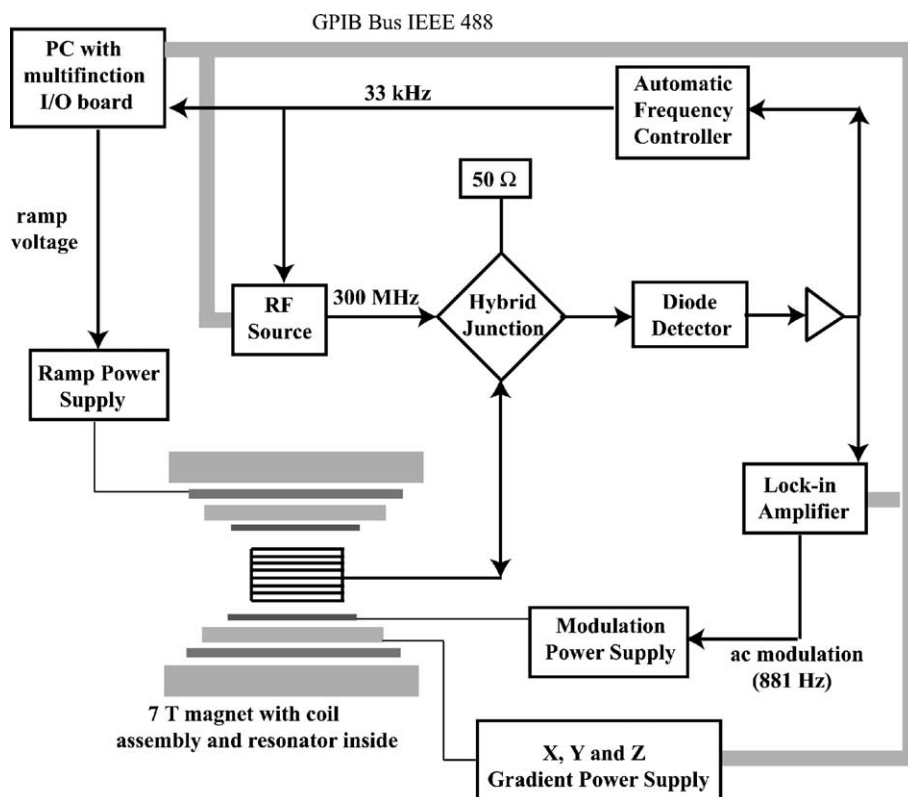


Fig. 1. Block diagram of the principal components in the CW-NMRI system.

excited plane of resonance is swept through the sample in the direction of the applied magnetic field gradient by the effect of the ramped offset magnetic field. A one-dimensional line profile is thus acquired, with each point in the profile representing the first derivative of the spin density of the sample projected onto the axis of the gradient. The gradient direction is rotated around the sample by the combination of three orthogonal gradients, and thus two- and three-dimensional images can be reconstructed from a series of projections using filtered back-projection.

2.2. Apparatus

The system is based around a 7 T, 183 mm diameter horizontal bore superconducting magnet (Oxford Instruments, UK). The gradient, offset, and modulation magnetic fields were produced using a custom-built water-cooled coil assembly (Laplacian, UK), which employed proprietary stream function designs to produce large gradient fields (maximum 400 mT/m for the Z and 300 mT/m for the X/Y gradients), together with two nested solenoids producing the ramped offset magnetic field (maximum ± 16 mT) and the superimposed audiofrequency modulation field (maximum ± 400 μ T). The modulation coil was placed on the inside layer of the assembly to minimise its inductance and hence current requirements, while the offset coil was placed on the outside to

minimise its coupling to the modulation coil. All magnetic fields were homogeneous to within 5% over a cylindrical volume of diameter 50 mm and length 70 mm. Commercial power supplies (Advance Hivolt, UK, Models AP90100 and AP5060) were used to power the gradient and offset coils (with maximum current capabilities of 90 and 50 A, respectively), while a home-built audiofrequency power supply was used to drive the modulation coil up to a maximum frequency of 15 kHz.

A range of birdcage resonators were built for studying the nuclei ^1H , ^{13}C , ^7Li , ^{23}Na , and ^{27}Al in a variety of materials (the same resonator was used to image ^{23}Na and ^{27}Al , tuned appropriately). In each case, it was necessary to avoid using materials in the construction of the resonators which contained the nuclei of interest, since these would be detected by the system and hence contribute a significant background signal. Thus, PTFE was used to provide structural support for the ^1H , ^7Li , and $^{23}\text{Na}/^{27}\text{Al}$ resonators, while the carbon-free glass ceramic Macor (Corning, USA) was used in the construction of the ^{13}C resonator. Specialised non-magnetic miniature chip ceramic capacitors (Tekelec, France) and variable capacitors (Voltronics, USA) were typically used throughout. The basic design of these birdcage resonators has been described previously [24], although some modifications were made here to suit specific applications. For instance, several ^1H resonators were built to separately optimise SNR, RF

uniformity, and sample size. The legs and end rings of the SNR-optimised resonator were made from 300- μm thick, 7 mm wide Cu foil, enabling an increase in surface area (hence decreased resistance) and a reduction in the number of solder joints compared to the original design, and with the end ring capacitors formed using appropriately shaped strips of CuFlon (Polyflon, USA), which is comprised of copper electroplated onto both sides of a 127- μm thick layer of PTFE serving as the capacitors' dielectric. The resonator optimised for RF uniformity differed from the original design in having 12 legs, while that optimised for sample size could accommodate a cylindrical sample with diameter 52 mm and length 100 mm. Similarly, the ^{13}C , ^7Li , and $^{23}\text{Na}/^{27}\text{Al}$ resonators were built to suit specific experimental requirements. All of the resonators were placed into identical RF shields made from copper sheets of thickness 0.5 mm and length 200 mm, with a resultant outer diameter of 81 mm. It was necessary to split the shields along their length to prevent the audiofrequency magnetic field modulation causing eddy currents, which would introduce acoustic interference and prevent the modulation from penetrating the shield to the sample. For the ^1H resonators, the split was closed to RF penetration by forming a capacitor comprising the shield itself and a strip of copper (width 18 mm, thickness 40 μm) bridging the gap, with a strip of PTFE (thickness 55 μm) in between acting as a dielectric. However, for the ^{13}C , ^7Li , and $^{23}\text{Na}/^{27}\text{Al}$ resonators, some RF leakage from the shields was detected due to the larger skin depths at these lower frequencies, and consequently Cu strips of width 80 mm and thickness 80 μm were used to bridge the gaps in these shields, at which point no RF leakage was detected. Measured Q -factors ranged from 132 for the $^{23}\text{Na}/^{27}\text{Al}$ resonator to 636 for the SNR-optimised ^1H resonator (with corresponding efficiencies of 1.7 and 14.4 $\mu\text{T}/\text{W}^{1/2}$, respectively).

Referring again to the block diagram in Fig. 1, the RF source, lock-in amplifier, and gradient power supplies were controlled through an IEEE 488 GPIB bus, while the ramp power supply was controlled by a voltage signal output through a multifunction board (National Instruments, USA, Model PCI 4060E) in the PC. This multifunction board was also used to control several home-built inverter units which were used to invert the current outputs from the unipolar gradient and ramp power supplies, and also to monitor the temperature at six points within the coil assembly. A home-built diode detector was used to convert the amplitude-modulated RF signal to a signal at the modulation frequency, which was subsequently amplified by a low-noise preamplifier. The resulting signal was split between the lock-in amplifier and a home-built automatic frequency controller (AFC). The feedback loop to the RF synthesiser via this AFC was used to compensate for drifts in the tuning of the various resonators during an experiment by adjusting

the RF source frequency (up to a maximum of ± 50 kHz for ^1H). Corresponding magnetic field offsets were performed by using the AFC feedback signal to modify the magnetic field ramp in software before it was sent to the ramp power supply. The sequence control, data acquisition, and data processing were carried out using software written in LabVIEW (National Instruments, USA), while image reconstruction was carried out using Interactive Data language (Research Systems, USA).

The resonator/shields were mechanically decoupled from the main gradient, offset, and modulation coil assembly by sliding them into a 1.6-m long fibreglass tube (1 mm wall thickness, inner diameter 83 mm) which was suspended within the coil assembly using two support frames built at either end of the magnet. The frames did not touch either the coil assembly or the magnet at any point and, furthermore, were mounted on vibration-dampening feet, in an effort to minimise microphonic effects due to vibrations originating in the modulation coil. A system for suspending and accurately positioning samples within the magnet was also incorporated into the support frame. Sample holders were made from PTFE to avoid background signal pick-up (except for the ^{13}C experiments, where Macor was used) and attached to a three-way translational stage on the positioning platform via a quartz rod.

2.3. Sample preparation and imaging parameters

2.3.1. ^{27}Al imaging experiments

A series of samples with water/cement (w/c) ratios of 0.3, 0.4, and 0.5 were made using a high temperature refractory cement trademarked Secar 80 (Lafarge Aluminates, France). Chemically, this calcium aluminate cement consists of approximately 80% Al_2O_3 and 17% CaO , together with trace amounts of other oxides. For 1-D profiling experiments, cylindrical samples measuring 44 mm in diameter and 60 mm in length were cast in cylindrical moulds and moist-cured in a wet environment (relative humidity $>95\%$) for 28 days to ensure a homogeneous cure. The samples were imaged using a magnetic field gradient of 50 mT/m, an RF power of 13 dBm, and a magnetic field modulation of 600 μT pk-pk. One hundred signal averages were typically used, resulting in an acquisition time of 10 min. For the 2-D imaging experiment, a sample with a rectangular cross-section was prepared (dimensions 60 \times 34 \times 32 mm) with a w/c ratio of 0.3 and again moist-cured in a wet environment (relative humidity $>95\%$) for 28 days. Two holes with diameters 6 mm and 8 mm were drilled through the sample. The following parameters were used to acquire the 2-D image: a magnetic field gradient of 80 mT/m, an RF power of 13 dBm, and a magnetic field modulation of 792 μT pk-pk, 150 signal averages, with 64 projections around the sample, resulting in a total acquisition time of 16 h.

2.3.2. ^{23}Na imaging experiments

A series of cylindrical samples (diameter 44 mm, length 60 mm) with w/c ratios of 0.3, 0.4, and 0.5 were made using “ordinary” Portland cement (OPC, made to European Standard EN197 but without interground limestone) and moist-cured in a wet environment (relative humidity >95%) for 28 days. The samples were subsequently dried in an oven at 105 °C for 3 days to remove all evaporable pore water, and immediately placed into a saturated solution of NaCl for 4 days, after which they were found to be completely saturated. In a second experiment, a cured, dried sample with a w/c ratio of 0.4 was sealed using Parafilm on all but one end face, which was exposed to a saturated solution of NaCl. The subsequent uptake of Na into the sample was followed by 1-D profiling experiments following different periods of soaking. In both experiments, the following parameters were used to acquire the profiles: a magnetic field gradient of 50 mT/m, an RF power of 13 dBm, and a magnetic field modulation of 792 μT pk-pk, 300 signal averages, resulting in a total acquisition time of 30 min.

2.3.3. ^{13}C imaging experiment

A cylindrical sample (diameter 44 mm, length 60 mm) with a w/c ratio of 0.35 was made from carbon-free Portland cement and cured in a desiccator with a relative humidity of 71%, and exposed to a constant flow of CO_2 gas for a period of 53 days, which was expected to fully carbonate the sample (this was verified by carrying out a phenolphthalein test on the sample after the imaging experiments were finished). The particular cement used for this experiment was chosen because it did not contain any calcite (i.e., calcium carbonate), which was verified by X-ray diffractometry. This was necessary in order to ensure that any measured ^{13}C signal derived from the carbonation reaction (i.e., from carbonated cement) rather than from admixtures within the cement. The sample was imaged using a magnetic field gradient of 30 mT/m, an RF power of 16 dBm, and a magnetic field modulation of 792 μT pk-pk. Because of the extremely low signal, it was necessary to use 8000 signal averages, resulting in a total acquisition time of 10 h.

For all images, the signal-to-noise ratio (SNR) was calculated as a ratio of the mean image intensity of a region of interest within the sample area compared to the standard deviation of the intensity.

3. Results and discussion

3.1. Imaging of solid phase of Al in cement

In addition to their binding properties, modern cements are characterised by a range of physical properties such as their rate of hardening, their resistance to ageing and temperature extremes, their rheological properties,

and their interaction with the fillers and additives typically used in the mix. These properties are determined by the mineralogical composition of the cement, which is generally tailored to suit the particular application. For example, for high temperature refractory applications, cements with high content of Al_2O_3 have been developed which combine rapid set and strength gain with high temperature resistance in subsequent service. Understanding the extent of the hydration reaction of the various components within these materials can aid in determining their long-term durability, and hence the ability to image the various phases in the cement is highly desirable.

One-dimensional profiling experiments carried out on the cylindrical Secar 80 cement samples as a function of different w/c ratios are presented in Fig. 2. It should be pointed out that we are imaging ^{27}Al nuclei predominantly in the chemically combined Al_2O_3 phase in the cement, although other solid Al-containing phases are also likely to contribute to the signal to a lesser extent. A T_2^* relaxation time of $49 \pm 1 \mu\text{s}$ was measured for the sample from the zero-gradient CW-NMR spectral linewidth. The decrease in the measured ^{27}Al signal of these samples as the w/c ratio is increased is clearly evident, demonstrating that the CW-NMRI system is sensitive to differences in the local cement concentration. It is interesting to note the slope to the right side (which corresponds to the top of the samples) for the w/c = 0.4 and 0.5 samples, which indicates that some degree of sedimentation may have taken place during the early stages of the cure. The raw 2-D image of the Secar 80 cement sample is presented in Fig. 3A. Both the 6 and 8 mm holes are resolved in this image, with some improvement observed following deconvolution of the image with its zero-gradient spectrum (Fig. 3B). This is further demonstrated in the horizontal line profiles taken through the mid-point of the two holes in each image, which are presented in Fig. 3C. The relatively poor spatial resolution is a direct result of the poor SNR of 19 which was measured in this image, which in itself reflects the poor quality factor ($Q = 132$) of the prototype ^{27}Al

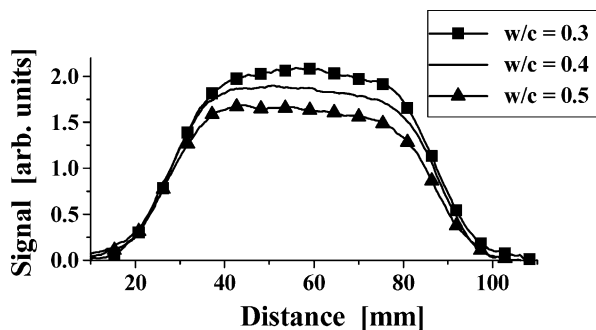


Fig. 2. 1-D profiles of the ^{27}Al content in cylindrical Secar 80 cement samples.

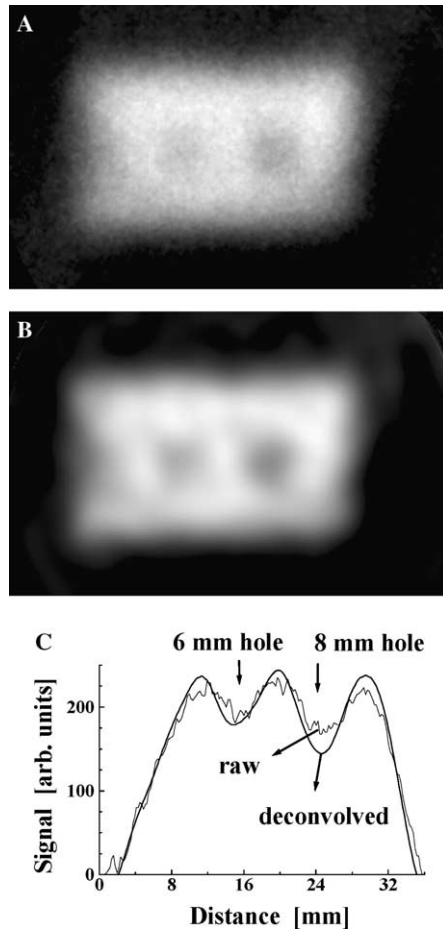


Fig. 3. 2-D image of the chemically combined ^{27}Al component in a Secar 80 cement sample, illustrating a 6 and 8 mm hole on the left and right side, respectively. (A) Raw image, (B) deconvolved image, and (C) a horizontal profile through the mid-point of the two holes for each image. A SNR of 19 was measured in the raw image.

resonator used to acquire the image. It is anticipated that a significant improvement in SNR can be achieved with an improved resonator design and a more rapid imaging sequence, which would allow for more signal averaging. Although a 1-D profile of ^{27}Al in an electrical insulating disk has been reported using STRAFI [26], to the authors' knowledge this is the first time that a solid aluminium phase in a cement material has been imaged in two dimensions using NMR techniques.

3.2. Imaging of salt penetration into cement

The ingress of salty water into concrete structures is of major economic concern due to the depassivation of embedded steel reinforcement resulting in corrosion, as well as direct damage caused by the effects of salt crystallisation. Bridges and coastal structures are particularly exposed to chloride attack via the ingress of salt water, while de-icing on concrete surfaces, e.g., bridge decks, via salt gritting can also cause significant problems. However, the movement of sodium and chloride

ions within cements is still not fully understood, while the reactions which may take place between them and the cement material itself or additives within the cement are equally vague. Indeed, interactions involving chlorides may retard their rate of penetration and thus provide some measure of long-term protection [32]. Several studies have been reported in the literature investigating the penetration of NaCl into porous building materials such as calcium silicate bricks [33] and OPC [30]. The current study aimed to explore the potential of the CW-NMRI technique in the study of salt water penetration into such materials.

One-dimensional profiling experiments carried out on the dried OPC cement samples, subsequently immersed in a saturated NaCl solution, are presented in Fig. 4. An increase in the ^{23}Na signal was measured as the w/c ratio was increased, due to the increased percolation with increased uptake of brine. In a second experiment, the rate of penetration of a brine solution into a dried and sealed sample as a function of soaking time is presented in Fig. 5. Spectral analyses indicated a two-phase component of the Na in the cement samples, exhibiting T_2^* values of ~ 90 and $\sim 180 \mu\text{s}$ with a relative intensity of approximately 1:2. The uncertainty in these values stems

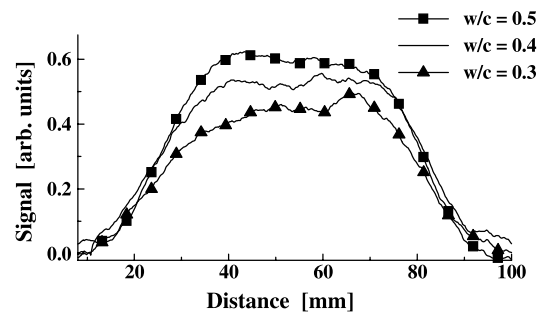


Fig. 4. Profiles of the ^{23}Na concentration in OPC samples with w/c ratios of 0.3, 0.4, and 0.5 soaked in brine, showing enhanced uptake for larger w/c values.

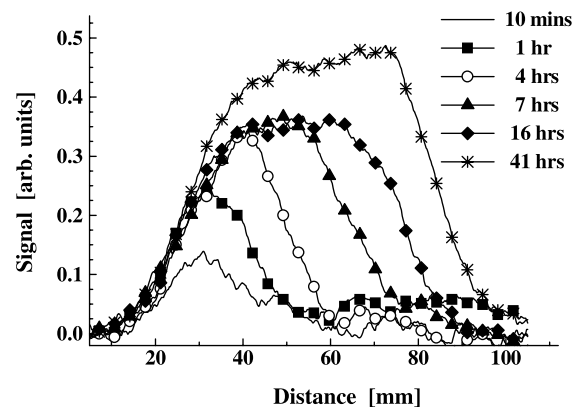


Fig. 5. Profiles of the ingress of ^{23}Na into an OPC sample with a w/c ratio of 0.4 as a function of soaking time.

from the difficulty in obtaining an accurate measure of the linewidths, since it was necessary to use an extremely small modulation amplitude of $\pm 20 \mu\text{T}$ to avoid modulation-broadening effects in the spectra, which resulted in an extremely small signal from the samples.

3.3. Imaging of the ^{13}C phase in carbonated cement

The carbonation of cement is another mechanism whereby the internal pH is lowered, thus depriving embedded steel of passivation. The process of carbonation typically takes place over a period of decades. Atmospheric CO_2 gas dissolved in water at the surface of the cement penetrates, initiating a chemical reaction involving the conversion of all the calcium-containing phases to, amongst other substances, calcium carbonate. For example, calcium hydroxide (which comprises ca. 25% of commonly used cements such as OPC) is converted to calcium carbonate and water. The resulting drop in pH levels within the cement exposes the steel reinforcement within the concrete structures to chemical attack by oxygen and chlorides [34]. However, the factors governing the rate of penetration of the carbonation front into cement materials are not fully understood and conventional techniques for measuring the rate of carbonation are destructive and hence prohibit repeat measurements on individual samples. The aim of the current study was therefore to investigate the feasibility of measuring a ^{13}C profile in a cement sample undergoing carbonation.

The ^{13}C content of the carbonated cement sample was imaged as detailed above. Without moving the sample, the ^{13}C resonator was removed and replaced with the large ^1H resonator, and further profiles were obtained using the same gradient strength and also a larger gradient of 300 mT/m to clearly delineate the edges of the sample. These profiles are presented in Fig. 6, in which the profiles have been scaled relative to each other for comparative purposes. The imaging of solid phases

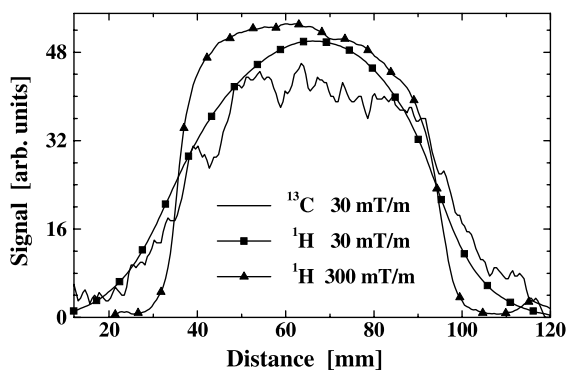


Fig. 6. 1-D profiles of the carbonated cement sample, showing ^{13}C and ^1H profiles measured at 30 mT/m and a high-resolution ^1H profile measured at 300 mT/m. The ^{13}C profile was multiplied by 500 for comparative purposes.

of ^{13}C in various rigid polymers has previously been reported [35,36]; however, to the authors' knowledge, this is the first time that a solid phase of ^{13}C in a carbonated cement sample has been imaged using NMR techniques. Due to the extremely low signal from this sample, it was not possible to get a measure of the T_2^* of the ^{13}C in this sample.

3.4. Imaging of ^7Li nuclei in powdered Li_2CO_3

A profile of the ^7Li nuclei in a cylindrical sample of powdered Li_2CO_3 (diameter 20 mm, length 35 mm) is presented in Fig. 7. The decrease in signal on the right side of the profile is an artefact caused by the severe signal saturation which occurred within this material, despite using an extremely low RF power of 1 mW (following baseline correction of the data, the right side of the profile was dropped onto the axis, artificially decreasing the signal in this part of the profile). The T_2^* of the ^7Li in this material was measured to be $55 \pm 6 \mu\text{s}$, with the relatively large error once again attributable to the small spectral signal obtained when using a small modulation amplitude.

3.5. 2-D and 3-D imaging of a rigid polymer

To investigate the maximum spatial resolution achievable with the system, a poly(methyl methacrylate) (PMMA) (Perspex/Plexiglas) resolution phantom was imaged using a gradient strength of 300 mT/m, an RF power of 10 dBm, and a modulation amplitude of 264 μT pk-pk, 6 signal averages and 180 projections around the sample, resulting in a total acquisition time of 1.8 h. Spectroscopic investigations carried out on this material using the CW-NMRI system indicated the presence of both crystalline and semi-amorphous phases (relative proportions 80:20), with T_2^* values of 16.3 and 26.9 μs , respectively. A schematic diagram of the phantom, together with the raw and deconvolved images, is presented in Fig. 8. A line profile taken through the 1 mm diameter hole in the raw image exhibited a decrease in signal intensity of 14%, indicating that it is

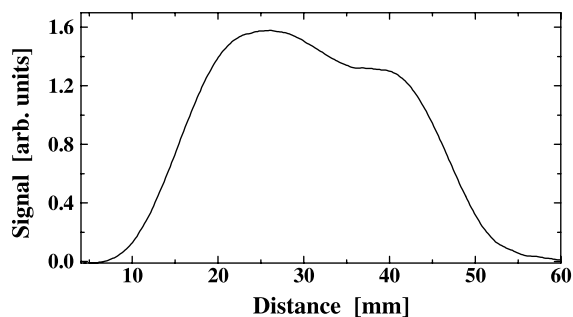


Fig. 7. 1-D profile of the ^7Li concentration measured in a sample of powdered Li_2CO_3 .

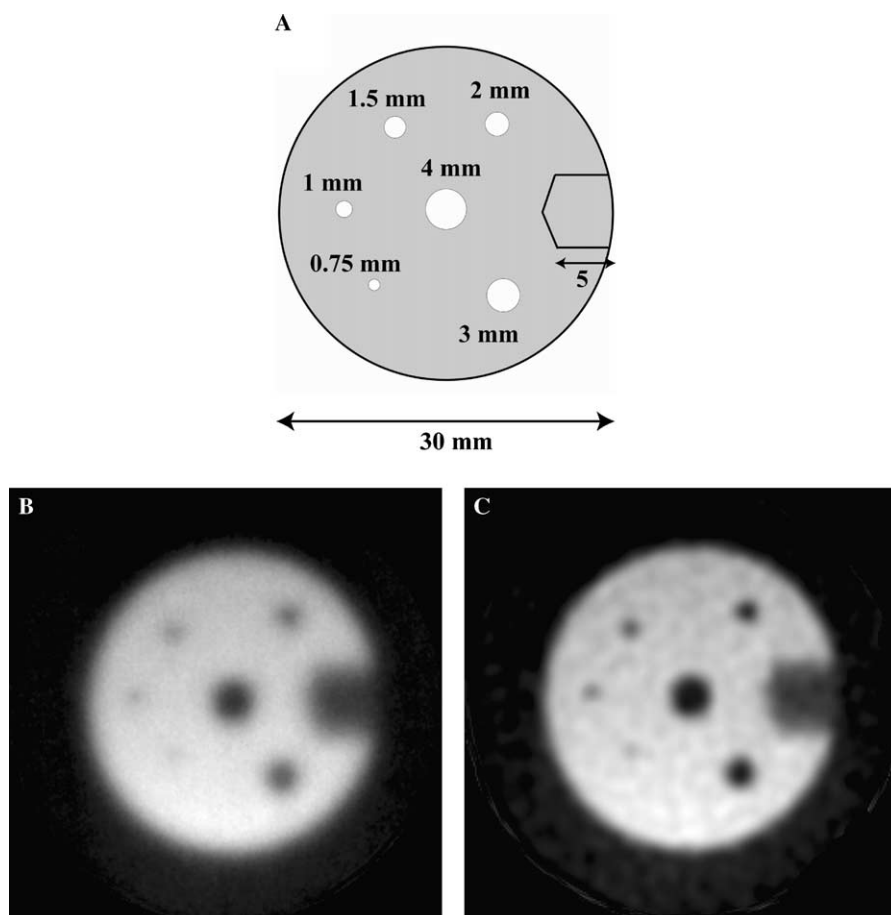


Fig. 8. (A) Schematic diagram of the PMMA resolution phantom, and (B) raw and (C) deconvolved images. A SNR of 82 was measured in the raw image.

on the verge of being resolved. This increased to 30% in the deconvolved image. The corresponding values for the 0.75 mm hole were 9 and 17% for the raw and deconvolved images, respectively. A SNR of 82 was measured in the raw image of Fig. 8B. By way of comparison, an early implementation of the SPRITE technique achieved a SNR of only 7 for a sample of PMMA following an image acquisition time of 15 min [37]. Using the same image acquisition time, a SNR of 34.1 was obtained with the CW-NMRI system.

To demonstrate the 3-D imaging capabilities of the new CW-NMRI system, the PMMA phantom shown schematically in Fig. 9A was constructed and imaged using a gradient strength of 300 mT/m, an RF power of 10 dBm, a modulation amplitude of 317 μ T pk-pk, and no signal averaging. The field of view was $(65 \text{ mm})^3$. In total, 5102 one-dimensional line projections were acquired at different angles around the sample, with the gradient direction stepped so as to produce isotropic resolution in the 3-D image. This resulted in an acquisition time of 8.5 h. The reconstructed dataset is shown as a series of 36 coronal slices through the imaged object in Fig. 9B, while 3-D surface

rendered views of the object are presented in Fig. 10. A SNR of 32.8 was measured in a 2-D slice positioned at roughly the mid-point through the 3-D dataset, with a resolution of approximately 1 mm. Comparisons with 3-D images of extremely short T_2 materials such as PMMA acquired using different solid imaging techniques are difficult due to the scarcity of published data in such materials. However, one such example was reported by Cory, who imaged a 6 mm diameter cylindrical PMMA phantom using a multi-pulse line-narrowing technique with the acquisition of 64×64 in-plane data elements [38]. A 1 mm hole in this phantom was well resolved in an axial slice through the image dataset, although the signal uniformity was relatively poor and no SNR value was quoted for this image. More recently, Halse et al. [31] acquired a ^{19}F 3-D image of a sample of PTFE exhibiting a slightly longer T_2^* of 35 μ s using the Centric SPRITE technique. With a $(64)^3$ image matrix and $(8 \text{ cm})^2 \times 20 \text{ cm}$ field of view, the resolution of this image was $(1.25 \text{ mm})^2 \times 3.125 \text{ mm}$, with a quoted value of S/N (ratio of the mean intensity of the signal and the mean intensity of the background noise) of 10.5 and acquisition time of 40 min. Although the 3-D

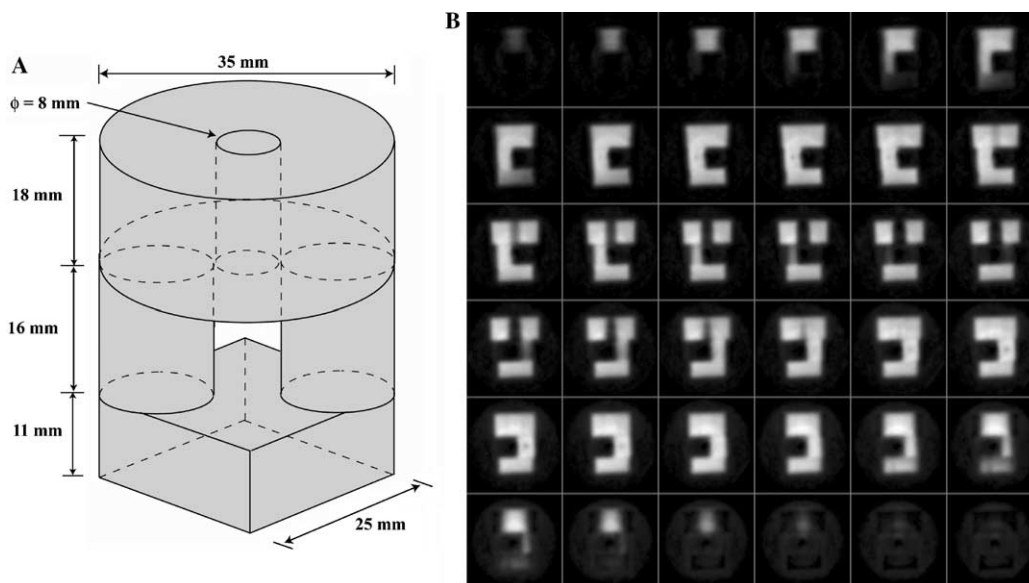


Fig. 9. (A) Schematic and (B) coronal slices through the reconstructed 3-D dataset of the PMMA phantom. A SNR of 32.8 was measured in a central slice.

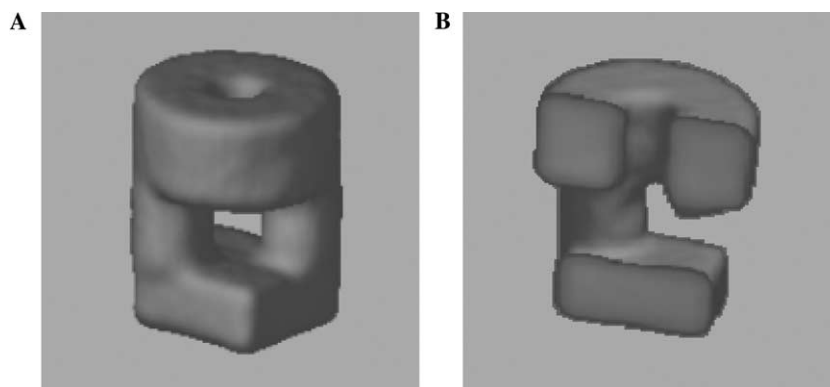


Fig. 10. 3-D surface rendered views of the PMMA phantom.

imaging performance of the CW-NMRI technique compares reasonably well with these other techniques, the relatively lengthy acquisition time is perhaps prohibitive for most practical imaging applications. Strategies for reducing the acquisition time are currently being investigated, and it is expected that significant time savings will be possible in future experiments.

The spatial resolution in the CW-NMRI images is ultimately limited by the available gradient strength and SNR, both of which could be improved following the continued development of the technique. For example, microphonic noise, which is the acoustic interference of the resonator from the solenoid coil producing the modulation field, was not fully eliminated in the new system but could be designed-out in a future system. Furthermore, considerable scope exists for refinement of both resonator design and image reconstruction techniques. If such improvements can

be implemented, the CW-NMRI technique may constitute a valuable addition to the arsenal of techniques currently available to study materials in the solid state.

4. Conclusions

The CW-NMRI technique has undergone a series of significant developments over recent years, to the point where complex, often heterogeneous materials with extremely short T_2 relaxation values can now be imaged routinely. This paper demonstrates, for the first time, continuous wave NMR imaging of nuclei other than ^1H , as well as the move into true 3-D imaging. The CW approach to imaging materials in the solid state has a number of advantages over more conventional pulsed MRI techniques. For example, the continual detection of the NMR signal eliminates any equipment

deadtime, while problems associated with the rapid switching of large gradients are avoided. The use of magnetic field modulation together with phase-sensitive detection renders CW-NMRI an inherently narrow-bandwidth detection technique. Perhaps more important, however, is the low RF power requirement of CW-NMRI, typically a factor of 10^3 – 10^5 lower than that required by other solid imaging techniques, thus making it practicable to envisage examination of full-size structural components which might prove problematic for conventional pulsed techniques due to excessive power disposition considerations. Although there exists scope for further improvements to the system, most notably in reducing imaging time (several strategies are currently being considered), the advantages expounded here together with the addition of 3-D, multi-nuclear imaging capabilities should lead to a broad range of applications for CW-NMRI across materials science.

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